

## 4-3 Coulomb's Law

Reading Assignment: *pp. 90-93*

Now, let' **determine** the electric field  $E(\vec{r})$  "produced" by charge density  $\rho_v(\vec{r})!$

Q:

A: HO: Coulomb's Law

Q:

A: HO: Coulomb's Law for Charge Distributions

# Coulomb's Law

Recall from **Coulomb's Law of Force** that a charge  $Q_2$  located at point  $\vec{r}_2$  applies a force  $\mathbf{F}_1$  on charge  $Q_1$  (located at point  $\vec{r}_1$ ):

$$\mathbf{F}_1 = \frac{1}{4\pi\epsilon_0} \frac{Q_1 Q_2}{R^2} \hat{a}_{21} = \frac{Q_1 Q_2}{4\pi\epsilon_0} \frac{\vec{r}_1 - \vec{r}_2}{|\vec{r}_1 - \vec{r}_2|^3}$$

Likewise, from the **Lorentz Force Law**, we know that the force  $\mathbf{F}_1$  on a charge  $Q_1$  located at point  $\vec{r}_1$  is attributed to an **electric field** located at  $\vec{r}_1$ :

$$\mathbf{F}_1 = Q_1 \mathbf{E}(\vec{r}_1) \quad \Rightarrow \quad \mathbf{E}(\vec{r}_1) = \frac{\mathbf{F}_1}{Q_1}$$

Inserting Coulomb's Law of Force into this equation, we get the electric field at location  $\vec{r}_1$ , generated by charge  $Q_2$  located at  $\vec{r}_2$ !

$$\mathbf{E}(\vec{r}) = \frac{\mathbf{F}_1}{Q_1} = \frac{Q_2}{4\pi\epsilon_0} \frac{\hat{a}_{21}}{R^2}$$

In general, we can say the electric field  $\mathbf{E}(\bar{r})$  at location  $\bar{r}$ , generated by a charge  $Q$  at point  $\bar{r}'$ , is:

$$\mathbf{E}(\bar{r}) = \frac{Q}{4\pi\epsilon_0} \frac{\hat{a}_R}{R^2} = \frac{Q}{4\pi\epsilon_0} \frac{\bar{r} - \bar{r}'}{|\bar{r} - \bar{r}'|^3}$$

This is **Coulomb's Law** !! It describes the electric field  $\mathbf{E}(\bar{r})$  at location  $\bar{r}$  that is created by a charge  $Q$  at location  $\bar{r}'$ .

Note that:

$$\hat{a}_R \doteq \frac{\bar{r} - \bar{r}'}{|\bar{r} - \bar{r}'|}$$

Therefore, if the charge  $Q$  is at the **origin** (i.e.,  $\bar{r}' = 0$ ), then:

$$\hat{a}_R = \frac{\bar{r}}{|\bar{r}|} = \hat{a}_r$$

Recall that the base vector  $\hat{a}_r$  always **points away** from the origin. In other words, a charge located at the origin creates an electric field vector that points in the direction of base vector  $\hat{a}_r$  (i.e., **away from the origin**) at all points  $\bar{r}$ !

Likewise, if the charge is at the origin, then:

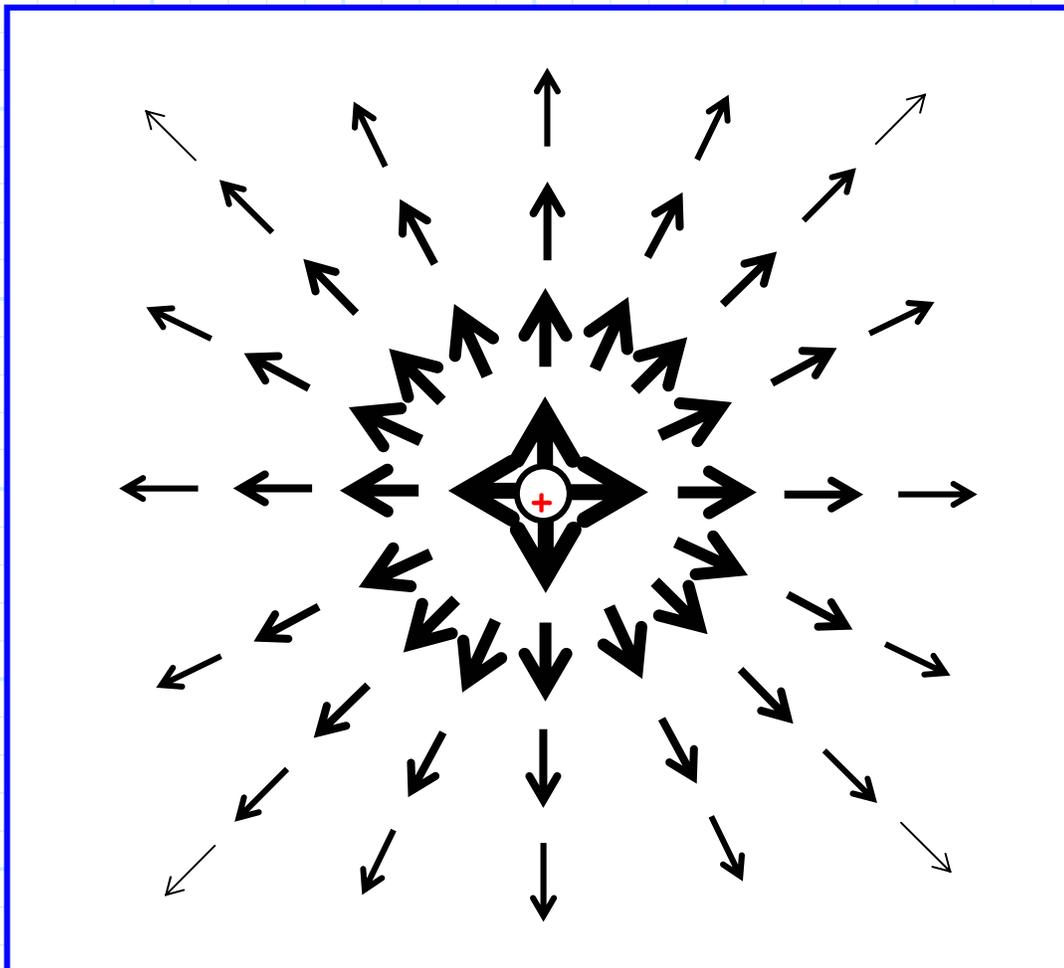
$$R = |\bar{r}| = r$$

In other words, the **magnitude** of the electric field vector is **proportional** to  $1/r^2$ . As a result, the magnitude of the electric field is dependent on its distance from the origin (i.e., distance from the charge). Therefore, if  $\vec{r}=0$ :

$$\mathbf{E}(\vec{r}) = \frac{Q}{4\pi\epsilon_0} \frac{\hat{a}_r}{r^2} = \frac{Q}{4\pi\epsilon_0} \frac{\vec{r}}{r^3}$$

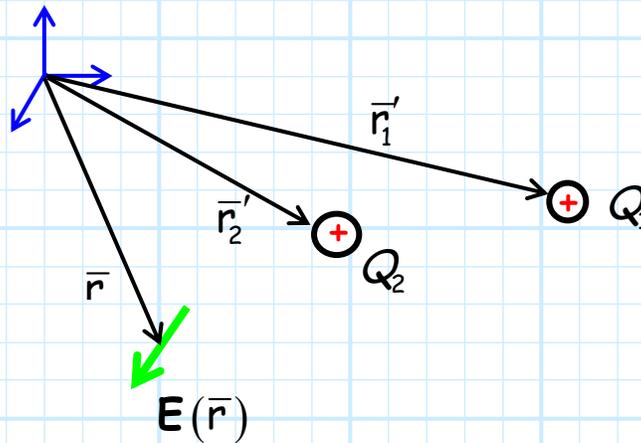
**Q:** What is the *curl* of  $\mathbf{E}(\vec{r})$  ??

**A:**  $\nabla \times \mathbf{E}(\vec{r}) =$



# Coulomb's Law for Charge Density

Consider the case where there are **multiple** point charges present. What is the resulting **electrostatic field**?



The electric field produced by the charges is simply the **vector** sum of the electric field produced by each (i.e., **superposition!**):

$$\mathbf{E}(\vec{r}) = \frac{Q_1}{4\pi\epsilon_0} \frac{\vec{r}-\vec{r}'_1}{|\vec{r}-\vec{r}'_1|^3} + \frac{Q_2}{4\pi\epsilon_0} \frac{\vec{r}-\vec{r}'_2}{|\vec{r}-\vec{r}'_2|^3}$$

Or, more generally, for  $N$  point charges:

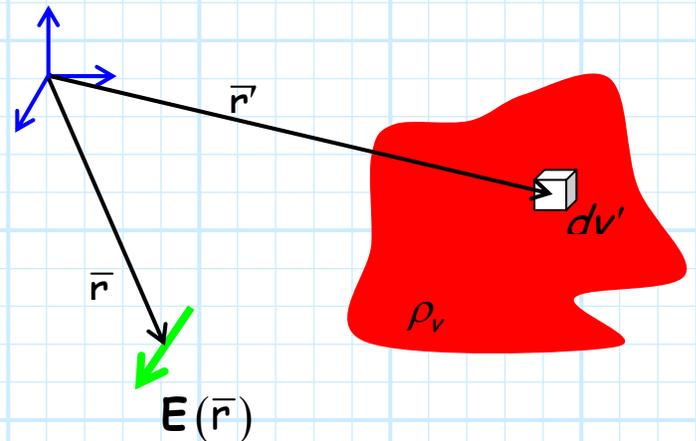
$$\mathbf{E}(\vec{r}) = \sum_{n=1}^N \frac{Q_n}{4\pi\epsilon_0} \frac{\vec{r}-\vec{r}'_n}{|\vec{r}-\vec{r}'_n|^3}$$

Consider now a volume  $V$  that is filled with a "cloud" of charge, described by **volume charge density**  $\rho_v(\bar{r})$ .

A very small differential volume  $dV$ , located at point  $\bar{r}'$ , will thus contain charge  $dQ = \rho_v(\bar{r}') dV'$ .

This differential charge produces an electric field at point  $\bar{r}$  equal to :

$$d\mathbf{E}(\bar{r}) = \frac{\rho_v(\bar{r}') dV'}{4\pi\epsilon_0} \frac{\bar{r} - \bar{r}'}{|\bar{r} - \bar{r}'|^3}$$



The **total** electric field at  $\bar{r}$  (i.e.,  $\mathbf{E}(\bar{r})$ ) is the summation (i.e., **integration**) of all the electric field vectors produced by all the little differential charges  $dQ$  that make up the charge cloud:

$$\mathbf{E}(\bar{r}) = \iiint_V \frac{\rho_v(\bar{r}')}{4\pi\epsilon_0} \frac{\bar{r} - \bar{r}'}{|\bar{r} - \bar{r}'|^3} dV'$$

**Note:** The variables of integration are the **primed** coordinates, representing the locations of the charges (i.e., **sources**).

Similarly, we can show that for **surface** charge:

$$\mathbf{E}(\bar{\mathbf{r}}) = \iint_S \frac{\rho_s(\bar{\mathbf{r}}')}{4\pi\epsilon_0} \frac{\bar{\mathbf{r}} - \bar{\mathbf{r}}'}{|\bar{\mathbf{r}} - \bar{\mathbf{r}}'|^3} ds'$$

And for **line** charge:

$$\mathbf{E}(\bar{\mathbf{r}}) = \int_C \frac{\rho_l(\bar{\mathbf{r}}')}{4\pi\epsilon_0} \frac{\bar{\mathbf{r}} - \bar{\mathbf{r}}'}{|\bar{\mathbf{r}} - \bar{\mathbf{r}}'|^3} d\ell'$$