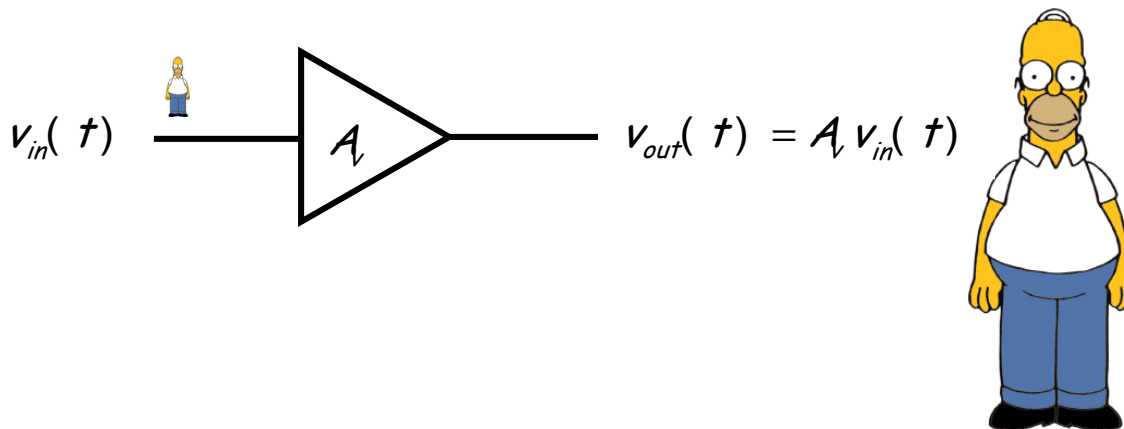


# Amplifiers

An **ideal** amplifier takes an input signal and reproduces it **exactly** at its output, only with a **larger** magnitude!



where  $A_v$  is the **voltage** gain of the amplifier.

Now, let's express this result using our **linear circuit theory** !

Recall, the output  $v_{out}(t)$  of a linear device can be determined by **convolving** its input  $v_{in}(t)$  with the device **impulse response**  $g(t)$ :

$$v_{out}(t) = \int_{-\infty}^t g(t-t') v_{in}(t') dt'$$

The impulse response for the **ideal** amplifier would therefore be:

$$g(t) = A \delta(t)$$

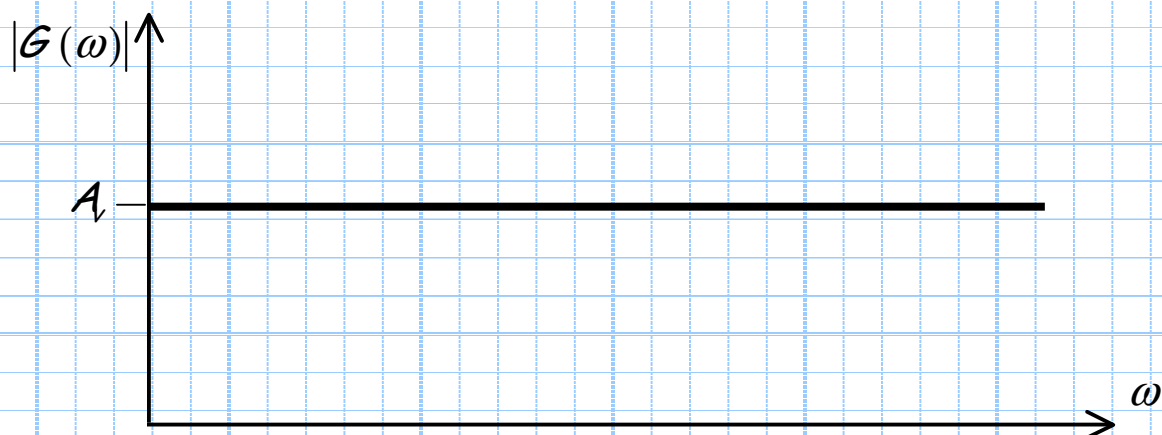
so that:

$$\begin{aligned}
 v_{out}(t) &= \int_{-\infty}^t g(t-t')v_{in}(t')dt' \\
 &= \int_{-\infty}^t A\delta(t-t')v_{in}(t')dt' \\
 &= Av_{in}(t)
 \end{aligned}$$

We can alternatively represent the ideal amplifier response in the **frequency domain**, by taking the **Fourier Transform** of the impulse response:

$$\begin{aligned}
 G(\omega) &= \int_{-\infty}^{\infty} g(t)e^{-j\omega t}dt \\
 &= \int_{-\infty}^{\infty} A\delta(t)e^{-j\omega t}dt \\
 &= A + j0
 \end{aligned}$$

This result, although simple, has an interesting interpretation. It means that the amplifier exhibits of magnitude gain of  $A_v$  for sinusoidal signals of **any** and **all** frequencies!



Moreover, the ideal amplifier does not alter the **relative phase** of the sinusoidal signal (i.e., no phase shift).



In other words, if:

$$v_{in}(t) = \cos(\omega t)$$

then at the output of the ideal amplifier we shall see:

$$\begin{aligned} v_{out}(t) &= |G(\omega)| \cos(\omega t + \angle G(\omega)) \\ &= A \cos(\omega t) \end{aligned}$$

BUT, there is one **big** problem with an ideal amplifier:



They are **impossible** to build !!

**Q:** *Why is that ??*

**A:** Two reasons:

- a) An ideal amplifier has **infinite** bandwidth.
- b) An ideal amplifier has **zero** delay.

**Not gonna happen !**

Let's look at this **first** problem first. The ideal amplifier impulse response  $g(t) = A \delta(t)$  means that the signal at the output occurs **instantaneously** with the signal at the input. This of course **cannot** happen, as it takes some small, but non-zero

amount of **time** for the signal to propagate through the amplifier. A more **realizable** amplifier impulse response is:

$$g(t) = A \delta(t - \tau)$$

resulting in an amplifier output of:

$$\begin{aligned} v_{out}(t) &= \int_{-\infty}^t g(t-t') v_{in}(t') dt' \\ &= \int_{-\infty}^t A \delta(t-\tau-t') v_{in}(t') dt' \\ &= A v_{in}(t-\tau) \end{aligned}$$

In other words, the output is both an amplified and **delayed** version of the input.

\* Note the delay does not **distort** the signal, as the output has the same form as the input.

\* Moreover, the delay for electronic devices such as amplifiers is **very small** in comparison to human time scales (i.e.,  $\tau \ll 1$  second).

\* Therefore, propagation delay  $\tau$  is **not** considered a **problem** for most amplifier applications.

BUT, the delay better be a **constant** with **frequency** (otherwise, signal **distortion** results)!

Let's examine what this delay means in the **frequency domain**.

Evaluating the Fourier Transform of this **modified** impulse response gives:

$$\begin{aligned}
 G(\omega) &= \int_{-\infty}^{\infty} g(t) e^{-j\omega t} dt \\
 &= \int_{-\infty}^{\infty} A \delta(t - \tau) e^{-j\omega t} dt \\
 &= A \cos(\omega\tau) + j A \sin(\omega\tau) \\
 &= A e^{j\omega\tau}
 \end{aligned}$$

We see that, as with the ideal amplifier, the magnitude  $|G(\omega)| = A$ . However, the relative **phase** is now:

$$\angle G(\omega) = \omega\tau$$

As a result, if  $v_{in}(t) = \cos(\omega t)$ , the output signal will be:

$$\begin{aligned}
 v_{out}(t) &= |G(\omega)| \cos(\omega t - \angle G(\omega)) \\
 &= A \cos(\omega t - \omega\tau)
 \end{aligned}$$

In other words, the output signal of a **real** amplifier is **phase shifted** with respect to the input.

Now, let's examine the **second problem** with the ideal amplifier. This problem is best discussed in the **frequency domain**.

We discovered that the **ideal** amplifier has a frequency response of  $|G(\omega)| = A_v$ . Note this means that the amplifier gain is  $A_v$  for **all** frequencies  $0 < \omega < \infty$  (D.C. to daylight!).

The bandwidth of the ideal amplifier is therefore **infinite**!

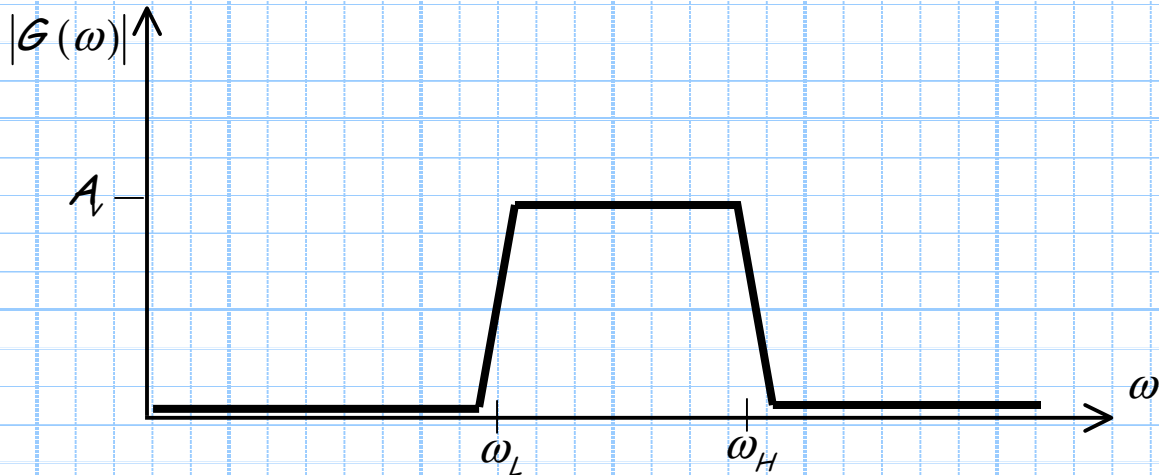
- \* Since every electronic device will exhibit **some** amount of inductance, capacitance, and resistance, every device will have a **finite** bandwidth.
- \* In other words, there will be frequencies where the device **does not work**!
- \* From the standpoint of an amplifier, "not working" means  $A_v < 1$  (i.e., **no gain**).
- \* Microwave/RF amplifiers will therefore have **finite** bandwidths.

There is a range of frequencies  $\omega$  between  $\omega_L$  and  $\omega_H$  where the gain will (approximately) be  $A_v$ . For frequencies outside this range, the gain will typically be small (i.e.  $A_v \approx 0$ ):

$$|G(\omega)| = \begin{cases} \approx A_v & \omega_L < \omega < \omega_H \\ \approx 0 & \omega < \omega_L, \omega > \omega_H \end{cases}$$

The width of this frequency range is called the amplifier bandwidth:

$$\begin{aligned} \text{Bandwidth} &\doteq \omega_H - \omega_L \quad (\text{radians/sec}) \\ &\doteq f_H - f_L \quad (\text{cycles/sec}) \end{aligned}$$



One result of having a **finite bandwidth** is that the amplifier impulse response is **not** an impulse function!

$$g(t) = \int_{-\infty}^{\infty} G(\omega) e^{+j\omega t} d\omega \neq A \delta(t - \tau)$$